

Conformal classes with explicit ambient metrics

Thomas Leistner



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Mathematical Sciences Institute, ANU
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[Joint work with Paweł Nurowski (Warsaw)]

Outline

Let $(M^n, [g])$ be a smooth manifold equipped with a conformal structure of semi-Riemannian metrics, $[g] := \{e^{2\varphi}g \mid \varphi \in C^\infty(M)\}$. Invariant descriptions:

- ① Fefferman-Graham ambient metric: Describe the conformal geometry of $(M^n, [g])$ in terms of the metric geometry of an *ambient space* $(\widetilde{M}^{n+2}, \widetilde{g})$ as in the flat case $S^n \simeq C^+/\mathbb{R}^+$ with $C^+ \subset \mathbb{R}^{1,n+1} =: \widetilde{M}$.
- ② Tractor bundle associated to the Cartan bundle with normal conformal Cartan connection.

Construct explicit ambient metrics for

- ① n -dimensional pp-waves
- ② Nurowski's conformal structures associated to $(2, 3, 5)$ -distributions with split G_2 Cartan connection \rightsquigarrow Holonomy split G_2 ambient metric

Construction of the ambient space

- A conformal class $[g]$ on M corresponds to an \mathbb{R}^+ -principle fibre bundle, the **cone**

$$\begin{array}{ccc} \bigcup_{x \in M} \{g_x \in \odot^2 T_x M \mid g \in [g]\} & =: C \cup \mathbb{R}^+, & \delta_t(g_x) := (t^2 g_x) \\ & \pi \downarrow & \\ & M & \end{array}$$

- **Tautological tensor** on C : $\mathbf{g}(U, V)|_{g_x} := g_x(d\pi(U), d\pi(V))$,
 - $\mathbf{g}(T, \cdot) = 0$ for the fundamental vector field T of δ .
 - of degree 2 w.r.t. the \mathbb{R}^+ -action, i.e. $\delta_t^* \mathbf{g} = t^2 \mathbf{g}$.
- Each $g \in [g]$ trivialises C ,

$$\begin{array}{ccc} \mathbb{R}^+ \times M & \simeq & C \\ (t, x) & \mapsto & (t^2 g_x, x) \end{array}$$

- Get the **ambient space** from the cone:

$$\widetilde{M} := C \times (-\varepsilon, \varepsilon)$$

The \mathbb{R}^+ action on C extends trivially to \widetilde{M} : $\delta_t(g_x, \rho) := (\delta_t g_x, \rho)$.

Ambient metric

Definition

Let $(M, [g])$ be smooth manifold with conformal class of signature (t, s) . An *ambient metric* for $(M, [g])$ is a smooth metric \tilde{g} on \tilde{M} such that

- 1 \tilde{g} is homogeneous of degree 2 w.r.t. the \mathbb{R}^+ -action δ .
- 2 $\iota^*\tilde{g} = \mathbf{g}$, for the inclusion $\iota : C = C \times \{0\} \subset \tilde{M}$.
- 3 $\text{Ric}(\tilde{g}) = 0$.

This implies that by fixing $g \in [g]$ trivialising C and a coordinate ρ such that $\tilde{M} \simeq \mathbb{R}^+ \times M \times (-\varepsilon, \varepsilon) \ni (t, x, \rho)$, we have

$$\tilde{g} = 2t d\rho dt + 2\rho dt^2 + t^2 g_\rho,$$

for a ρ -dependent family of metrics g_ρ with $g_0 = g$ and subject to the condition $\text{Ric}(\tilde{g}) = 0$.

The Fefferman-Graham ambient metric

Theorem (C. Fefferman & C.R. Graham '85, '07)

- If $n := \dim M$ is odd, then
 - ① there exists a formal power series solution g_ρ to $\text{Ric}(\bar{g}) = 0$,
 - ② this solution is unique up to \mathbb{R}^+ invariant diffeomorphisms fixing $C \subset \bar{M}$,
 - ③ if $[g]$ is analytic, then there exists an ambient metric.
- If n is even, then there exists a formal power series solution to $\text{Ric}(\bar{g}) = O(\rho^{\frac{n-2}{2}})$ which is uniquely determined up to terms of order $\frac{n}{2}$ in ρ . Furthermore, there exists a conformally invariant tensor $O \in \Gamma(\odot^2 TM)$, the obstruction tensor, such that $O \equiv 0 \iff \text{Ric}(\bar{g}) = 0$ to infinite order.

O is trace and divergence free, conf. invariant of weight $(2 - n)$, given by

$$O = \Delta_g^{n/2-2} (\Delta_g P - \nabla^2 \text{tr}(P)) + \text{lower order terms},$$

where $P = \frac{1}{n-2} \left(\text{Ric} - \frac{\text{scal}}{2(n-1)} g \right)$ is the Schouten tensor and Δ_g is the Laplacian of $g \in [g]$.

First terms for g_ρ

Expanding g_ρ as a power series we get

$$\tilde{g} = 2t d\rho dt + 2\rho dt^2 + t^2 \left(g + \sum_{k=1}^{\infty} \rho^k \mu_k \right).$$

with

$$\begin{aligned} (\mu_1)_{ab} &= 2P_{ab} \\ (n-4)(\mu_2)_{ab} &= -B_{ab} + (n-4)P_a{}^c P_{bc} \\ 3(n-4)(n-6)(\mu_3)_{ab} &= \Delta_g B_{ab} + \dots \end{aligned}$$

where $B_{ab} = \nabla_c C_{ab}{}^c - P_{cd} W_{ab}{}^{cd}$ is the Bach tensor, W_{abcd} is the Weyl tensor, $C_{abc} := \nabla_c P_{ab} - \nabla_b P_{ac}$ is the Cotton tensor.

Ambient metrics for Einstein classes

$[g]$ contains a Einstein metric $\implies \tilde{g}$ admits a parallel vector field:
 If $g_\Lambda \in [g]$ is Einstein with $P = \Lambda g$. Then

$$\tilde{g} = 2d(\rho t)dt + (t^2 + 2t^2\rho\Lambda + t^2\rho^2\Lambda^2)g_\Lambda$$

$$\stackrel{-u=t\rho}{=} -2dudt + (t^2 - 2\Lambda ut + \Lambda^2 u^2)g_\Lambda$$

- $\Lambda \neq 0$: $\tilde{g} \stackrel{\substack{r=t-\Lambda u \\ s=t+\Lambda u}}{=} -\frac{1}{2\Lambda}ds^2 + \underbrace{\frac{1}{2\Lambda}dr^2 + r^2g_\Lambda}_{\text{cone metric}}$

- $\Lambda = 0$: $\tilde{g} = -dudt + t^2g_\Lambda$.

Further examples of explicit ambient metrics with truncated ambient metric:

- Products of Einstein metrics with related Einstein constants [R. Gover & F. Leitner '06]

pp-waves

An n -dimensional Lorentzian (pseudo-Riemannian) manifold (M, g) is a *pp-wave* (“plane fronted with parallel rays”) \iff

- \exists parallel null vector field K ,
- $R(U, V) : K^\perp \rightarrow \mathbb{R}K$ for all $U, V \in TM$

$\iff (M, g)$ has abelian holonomy contained in $\mathbb{R}^{n-2} \subset \text{SO}(n-2) \ltimes \mathbb{R}^{n-2}$

$\iff (M, g)$ admits coordinates $(x^1, \dots, x^{n-2}, u, r)$ such that

$$g = \sum_{i=1}^{n-2} (dx^i)^2 + 2du (dr + hdu), \quad \text{with } h = h(x^1, \dots, x^{n-2}, u)$$

All scalar invariants vanish, in particular $\text{scal} = 0$ and

$$P = \frac{1}{n-2} \text{Ric} = -\frac{1}{n-2} \Delta h du^2 \quad \text{and} \quad B = -\frac{1}{n-2} \Delta^2 h du^2$$

where $\Delta h = \sum_{i=1}^{n-2} \partial_i^2 h$.

Note: Einstein $\iff \text{Ric} = 0 \iff \Delta h = 0$.

The ambient metric of a pp-wave

Let $(M, g = \sum_{i=1}^{n-2} (dx^i)^2 + 2dudr + 2hd u^2)$ be a pp-wave.

Ansatz for ambient metric

$$\begin{aligned}\tilde{g} &= 2d(\rho t)dt + t^2(g + Hd u^2) \\ &= 2d(\rho t)dt + t^2\left(2dudr + 2(h + H)du^2 + \sum_{i=1}^{n-2} (dx^i)^2\right),\end{aligned}$$

where $H = H(\rho, x^i, u)$, and $H(\rho, x^i, u)|_{\rho=0} = 0$. Then

$$\text{Ric}(\tilde{g}) = \left((2-n)\partial_\rho H + 2\rho\partial_\rho^2 H - \Delta H - \Delta h \right) du^2,$$

I.e., \tilde{g} is ambient metric for $[g]$ if

$$\left((2-n)\partial_\rho + 2\rho\partial_\rho^2 - \Delta \right) (H) = \Delta h. \quad (1)$$

Solution to (1) is given by

$$H := \sum_{k=1}^{\infty} \frac{\rho^k \Delta^k h}{k! \prod_{i=1}^k (2i - n)}$$

Theorem (Nurowski & L. '10)

Let $g = \sum_{i=1}^{n-2} (dx^i)^2 + 2du (dr + hdu)$ be a pp-wave metric. If n is odd, then the ambient metric for \bar{g} for the conformal class $[g]$ is given by

$$\bar{g} = 2d(\rho t)dt + t^2 g + t^2 \left(\sum_{k=1}^{\infty} \frac{\rho^k \Delta^k h}{k! \prod_{i=1}^k (2i - n)} \right) du^2.$$

If n is even, the obstruction tensor is given as $O = \Delta^{\frac{n}{2}} h du^2$. If $O = 0$, then

$$\bar{g} = 2d(\rho t)dt + t^2 g + t^2 \left(\sum_{k=1}^{\frac{n-2}{2}} \frac{\rho^k \Delta^k h}{k! \prod_{i=1}^k (2i - n)} \right) du^2. \quad (2)$$

is an ambient metric.

Note that even if $O = \Delta^{\frac{n}{2}} h du^2 = 0$ the ambient metric is not unique since we can add terms of the form

$$\rho^{\frac{n}{2}} \Delta^{\frac{n-2}{2}} h du^2$$

in (2) and still get a Ricci flat ambient metric.

Bach flat pp-waves in dimension 4

In dimension $n = 4$ the obstruction to the existence of the ambient metric for a pp-wave

$$g = dzd\bar{z} + 2dudr + 2hdu^2$$

is given by the Bach tensor $B = -\frac{1}{n-2}\Delta^2 hdu^2$.

Proposition

A 4-dimensional pp-wave is Bach flat \iff

$$h = z\bar{\alpha} + \bar{z}\alpha + \beta + \bar{\beta} \quad \text{for } \alpha = \alpha(z, u) \text{ and } \beta = \beta(z, u) \text{ holom. in } z$$

Furthermore, a 4-dimensional Bach flat pp-wave satisfies

$$g \text{ is conformally flat } \iff g \text{ is conformally Einstein } \iff \partial_z^2 \alpha = 0.$$

\rightsquigarrow Examples for which the ambient metric truncates but which are not conformally Einstein.

Normal conformal tractor bundle

Let $(M, [g])$ be a conformal manifold of signature (p, q) and set $G := SO^0(p+1, q+1)$, $P = \text{Stab}_G(L)$ and $\mathfrak{g} := \mathfrak{so}(p+1, q+1)$:

- $\omega \in T\mathcal{P} \rightarrow \mathfrak{g}$ be the normal conformal Cartan connection on the Cartan bundle \mathcal{P} with structure group P .
- $\bar{\omega} : T\bar{\mathcal{P}} := T(\mathcal{P} \times_P G) \rightarrow \mathfrak{g}$ extension to a G -bundle connection, and
- $\bar{\nabla}$ the corresponding covariant derivative on the tractor bundle $\mathbb{T} = \bar{\mathcal{P}} \times_G \mathbb{R}^{p+1, q+1}$.

Every $g \in [g]$ splits $\mathbb{T} = \underline{\mathbb{R}} \oplus TM \oplus \underline{\mathbb{R}}$ with

$$\bar{g} = \begin{pmatrix} 0 & 0 & 1 \\ 0 & g & 0 \\ 1 & 0 & 0 \end{pmatrix} \quad \text{and} \quad \bar{\nabla}_X \begin{pmatrix} \tau \\ Y \\ \sigma \end{pmatrix} = \begin{pmatrix} d\tau(X) - P(X, Y) \\ \nabla_X Y + \tau X + \sigma(X \lrcorner P)^\# \\ d\sigma(X) - g(X, Y) \end{pmatrix}.$$

Relation to the ambient metric

Computing the LC connection $\widetilde{\nabla}$ of \widetilde{g} along $M \subset \widetilde{M}$ yields that the vector bundle isomorphism

$$\Phi : T\widetilde{M}|_M \ni \sigma \frac{\partial}{\partial \rho} + \tau \frac{\partial}{\partial t} + Y \mapsto \begin{pmatrix} \tau \\ Y \\ \sigma \end{pmatrix} \in \mathbb{T}$$

is affine for $\widetilde{\nabla}$ and $\overline{\nabla}$, i.e. $\Phi(\widetilde{\nabla}_X U) = \overline{\nabla}_X(\Phi(U))$. In particular, $\text{Hol}_x(\mathbb{T}, \overline{\nabla}) = \text{Hol}_x(T\widetilde{M}|_M, \widetilde{\nabla})$ and hence

$$\text{Hol}_x(\mathbb{T}, \overline{\nabla}) \subset \text{Hol}_x(T\widetilde{M}, \widetilde{\nabla}). \quad (3)$$

- Extension of parallel tractors to \widetilde{M} yields = in (3) at least when n is odd: Graham & Willse [last week].

Parallel objects for the tractor connection

Let $(M, [g])$ be a conformal manifold.

“ $[g]$ is conformal ...” :=

“on an open dense subset of M there is a metric g in $[g]$ such that g is ...”

- ① $\bar{\nabla}$ admits a parallel line bundle $\iff [g]$ is conformal Einstein.
- ② $\bar{\nabla}$ admits a parallel plane bundle of totally null planes $\iff [g]$ is conformal *aligned pure radiation*. g is aligned pure radiation metric if
 - \exists parallel null line bundle \mathcal{K} ,
 - $\mathcal{K}^\perp \lrcorner Ric = 0$.
- ③ $[g]$ is a conformal pp-wave $\iff \bar{\nabla}$ admits a parallel totally null plane subbundle \mathcal{H} such that \bar{R} satisfies

$$\bar{R}_{abCD} X^C Y^D = 0,$$

for all X^C and Y^D orthogonal to \mathcal{H} . [2,3: L '06, Nurowski & L '11].

Replace tractor connection $\bar{\nabla}$ by ambient connection $\tilde{\nabla}$ we get \implies in the above statements.

Nurowski's conformal structures

For a function F on $\mathbb{R}^5 = \{(x, y, p, q, z)\}$ with $F_{qq} \neq 0$ consider the $(2, 3, 5)$ -distribution \mathcal{D}_F spanned by

$$\partial_q \text{ and } \partial_x + p\partial_y + q\partial_p + F\partial_z.$$

\mathcal{D}_F defines a conformal class $[g_F]$ on $\mathbb{R}^{3,2}$ such that $\mathcal{D}_F^\perp = \text{span}(\mathcal{D}_F, [\mathcal{D}_F, \mathcal{D}_F])$ and whose normal conformal Cartan connection reduces to split G_2 .

Nurowski observed: For F of the form

$$F = q^2 + \sum_{i=1}^6 a_i p^i + bz \tag{4}$$

the ambient metric for $[g_F]$ is of the form

$$\tilde{g}_F = -2dtdu + t^2 g_F - 2utP - u^2 B, \tag{5}$$

where P and B are the Schouten and Bach tensors of g_F and $u = -t\rho$.

Properties of $[g_F]$

- ① $F = F(q) \implies [g_F]$ contains Ricci-flat metric $\implies \text{Hol}(\tilde{g}) \subset \text{Stab}_{G_2}(\text{null vector})$.
- ② If at least one of a_4 , a_5 , or a_6 is not zero, then the conformal class $[g_F]$ is not conformal Cotton flat and thus, not conformal Einstein.
- ③ For $a_4 = a_5 = a_6 = 0$ and $a_3 \neq 0$ there is a unique Cotton flat metric in $[g_F]$, but $[g_F]$ is not conformal Einstein.
- ④ If at least one of a_3 , a_4 , a_5 , a_6 not equal to zero. Then $[g_F]$ is not conformal aligned pure radiation.

Proposition

For F as in (4) the ambient metric \tilde{g}_F of $[g_F]$ admits a parallel spinor of non-zero length. Hence, $\text{Hol}(\tilde{g}_F) \subset G_2$.

Ambient holonomy G_2

Theorem (Nurowski & L '11)

If $\text{Hol}(\widetilde{g}_F) \neq G_2$, then $[g_F]$ is conformal Einstein or conformal aligned pure radiation. In particular, if at least one of a_3, a_4, a_5 or a_6 in F is not zero then $\text{Hol}(\widetilde{g}_F) = G_2$.

Proof: Berger's list in signature $(4, 3)$: G_2 and $SO^0(4, 3)$. Have to check that $\text{Hol}(\widetilde{g}_F)$ has no invariant subspace. Let $V \subset \mathbb{R}^{4,3}$ be invariant.

- V non-degenerate. De Rham $\Rightarrow \widetilde{g}_F$ is locally a product with both factors Ricci-flat. One factor has $\dim \leq 3 \Rightarrow \text{flat} \Rightarrow \text{conformal Ricci flat}$
- $V \cap V =: W \neq 0$ totally null.
 - $\dim(W) = 1 \Rightarrow \text{conformal Ricci flat.}$
 - $\dim(W) = 2 \Rightarrow \text{conformal aligned pure radiation.}$
 - $\dim(W) = 3$ maximal null space $\Rightarrow \exists$ invariant null line of spinors. Then

Non-null spinor ψ + null line of spinors $\mathbb{R}\varphi \Rightarrow$ null line of vectors $\mathbb{R}V_{\psi,\varphi}$

via $\widetilde{g}_F(V_{\psi,\varphi}, X) = \langle X \cdot \psi, \varphi \rangle$. Hence, $[g_F]$ is conformal Ricci flat. \square