

Acoustic metamaterial chains involving inertial amplification: *Linear case*

Luke Bennetts

University of Adelaide



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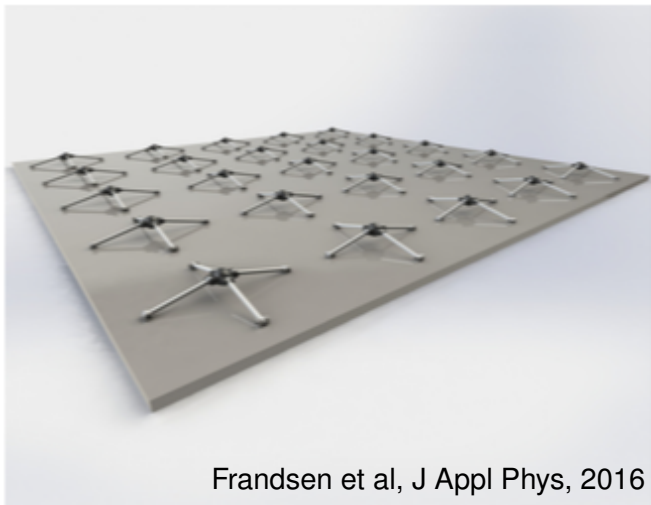
Low-frequency noise



Problem

- Low-frequency noise travels farthest.
- Most difficult to isolate (or damp).
- AMMs are designed to isolate low-frequency noise. ✓
- Typically require heavy elements. ✗

AMMs using lightweight attachments



Frandsen et al, J Appl Phys, 2016

- Idea from J-dampers in F1 racing car suspension systems.

AMM as 1D chain

Metamaterials — according to Wikipedia

"A material engineered to have a property that is not found in nature"

1D chain

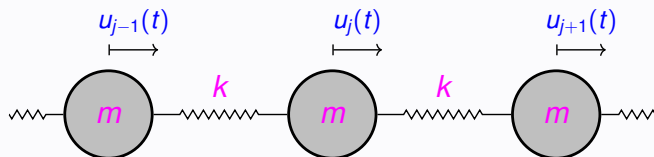


where something *"meta"* happens in the boxes.

Simple starting point for analysis

Monatomic chain: Not an AMM

Infinite chain with parameters $m = \text{mass}$ and $k = \text{spring}$



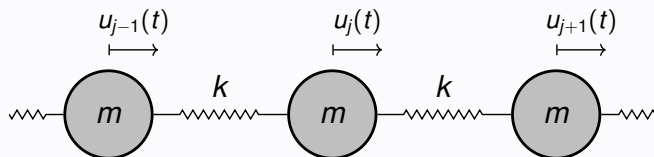
Equations of motion

$$m \ddot{u}_j + k (2 u_j - u_{j-1} - u_{j+1}) = 0 \quad \text{for } j \in \mathbb{Z}$$

Simple starting point for analysis

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Infinite chain with parameters m = mass and k = spring



Equations of motion

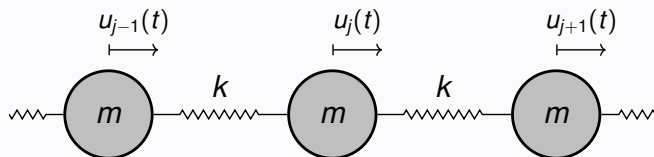
$$-\omega^2 m \hat{u}_j + k (2 \hat{u}_j - \hat{u}_{j-1} - \hat{u}_{j+1}) = 0$$

- Fourier transform, i.e. write $u_j(t) = \hat{u}_j \exp(-i \omega t)$.

Simple starting point for analysis

Monatomic chain: Not an AMM

Infinite chain with parameters m = mass and k = spring



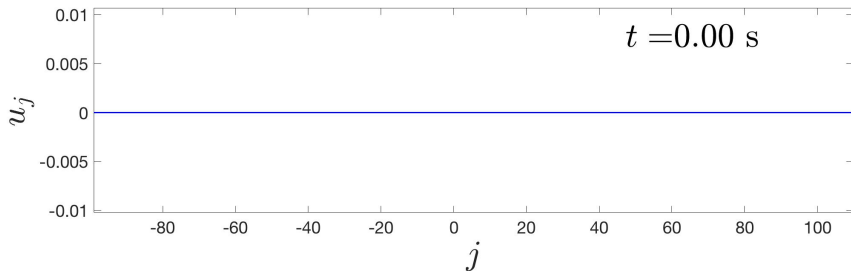
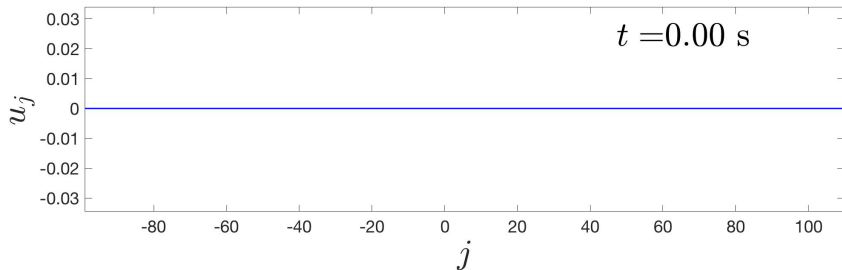
Dispersion relation

$$m \omega^2 = 2 k (1 - \cos q)$$

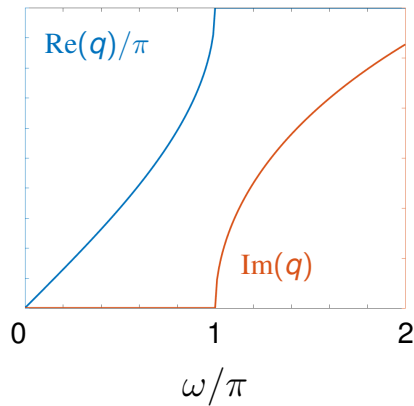
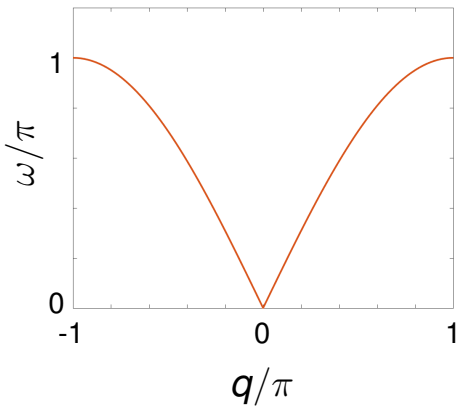
- Fourier transform, i.e. write $u_j(t) = \hat{u}_j \exp(-i \omega t)$.
- Apply Bloch–Floquet condition (quasi-periodicity)

$$\hat{u}_j = U \exp(i q j) \quad \text{where} \quad q = \text{wavenumber} \in (-\pi, \pi].$$

$m = 1$; $k = \pi^2/4$; $\omega/\pi = 0.6$ (top), 1.2 (bottom)

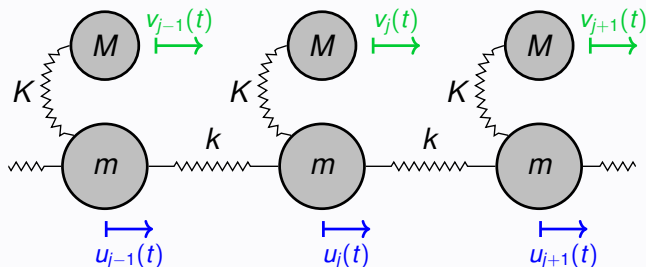


Band diagrams: $m = 1$; $k = \pi^2/4$



Turn it into a metamaterial

Standard AMM: Mass-in-mass chain



Equations of motion

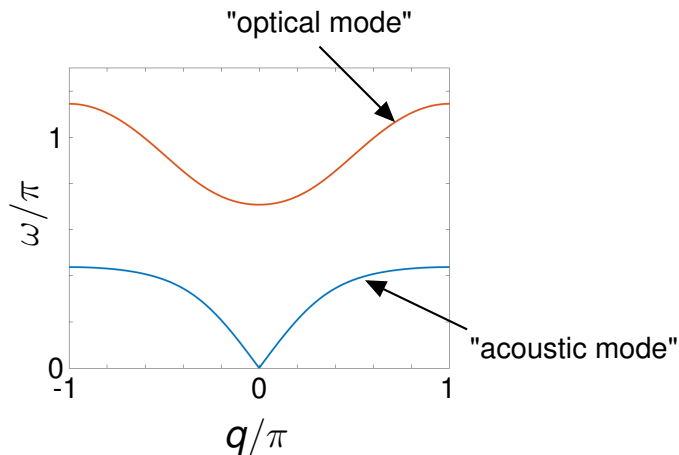
$$m \ddot{u}_j + k (2 u_j - u_{j-1} - u_{j+1}) + K(u_j - v_j) = 0$$

$$M \ddot{v}_j + K (v_j - u_j) = 0$$

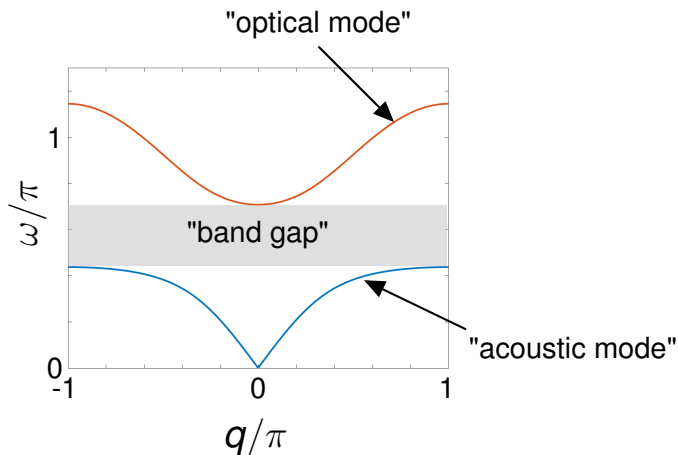
Dispersion relation

$$m M \omega^4 - \{K (m + M) + 2 k (1 - \cos q)\} \omega^2 + 2 (1 - \cos q) = 0$$

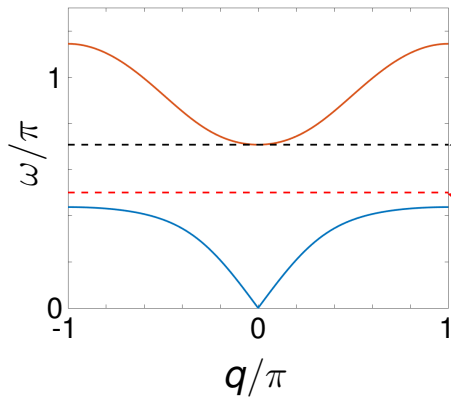
Band diagrams: $M = m = 1$; $K = k = \pi^2/4$



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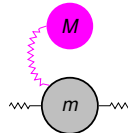
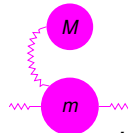


Band diagrams: $M = m = 1$; $K = k = \pi^2/4$



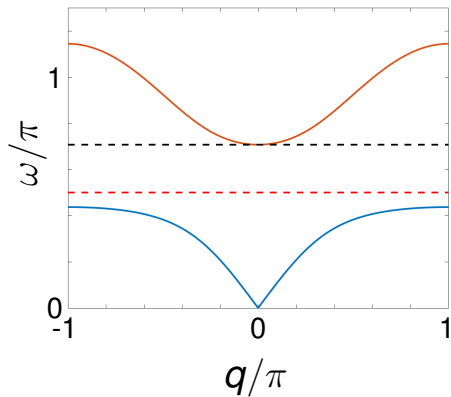
resonance: $\omega^2 = \omega_0^2 \equiv \frac{K(m+M)}{mM}$

anti-resonance: $\omega^2 = \omega_*^2 \equiv \frac{K}{M}$



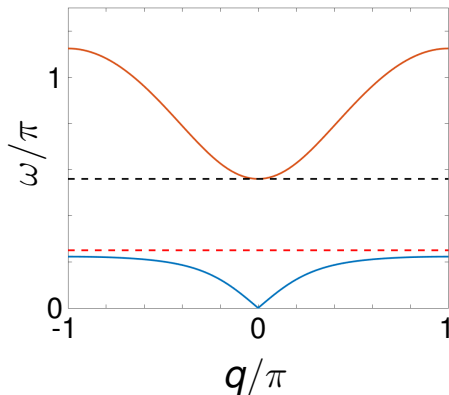
Band diagrams: $m = 1$; $K = k = \pi^2/4$

$$M = m$$



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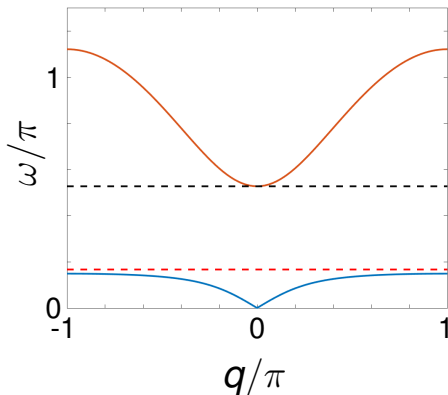
$$M = 4m$$



Lower resonances
= lower band gap

Band diagrams: $m = 1$; $K = k = \pi^2/4$

$$M = 9m$$



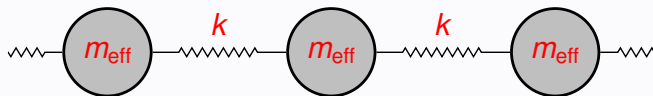
Lower resonances
= lower band gap

But mass increases

Why is mass-in-mass chain an AMM?

"... a property that is not found in nature"

Mass-in-mass chain \mapsto "effective" monatomic chain



with dispersion relation (from earlier)

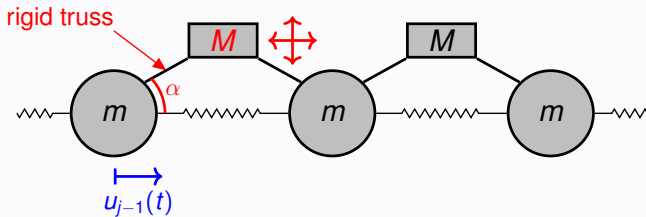
$$\frac{m_{\text{eff}} \omega^2}{k} = 2(1 - \cos q) \in [0, 4] \quad \text{for } q \in \mathbb{R}$$

where

$$m_{\text{eff}} = m + \frac{M \omega_*^2}{\omega_*^2 - \omega^2} : \quad |\omega_*| \gg 1 \quad \text{for } \omega \sim \omega_*$$

AMM with lightweight attachments (+ geometry)

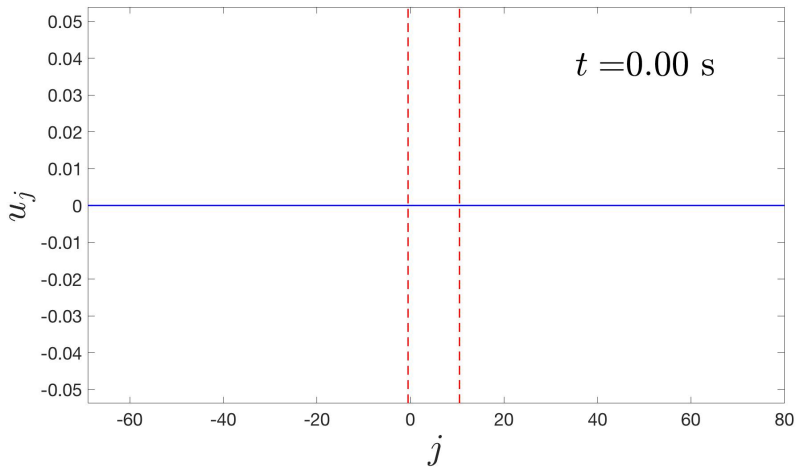
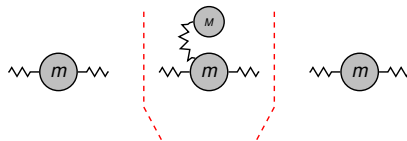
Simplest variant (Yilmaz & Hulbert, Phys. Lett. A, 2010)



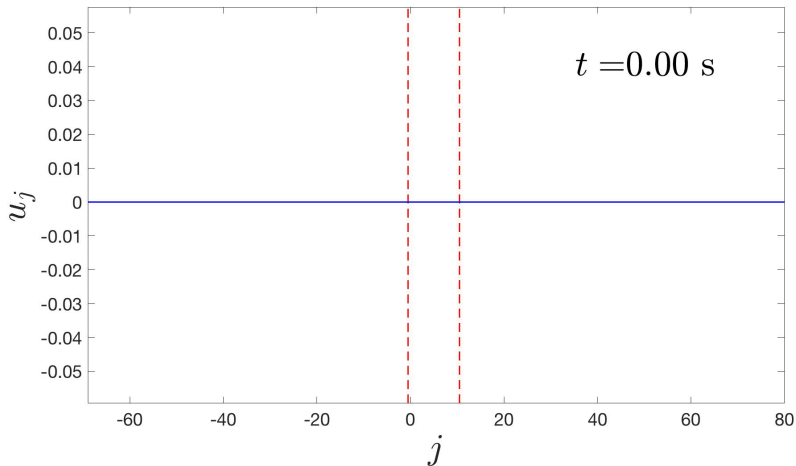
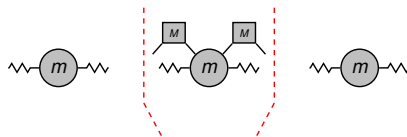
Equation of motion

$$\left(m + \frac{1}{4} M (1 + \cot^2 \alpha) \right) \ddot{u}_j + \frac{1}{4} M (1 - \cot^2 \alpha) (\ddot{u}_{j-1} + \ddot{u}_{j+1}) + k (2 u_j - u_{j-1} - u_{j+1}) = 0$$

Mass-in-mass: $M = m/10$; $\omega = 6\pi/10$



Yilmaz & Hulbert: $M = m/9$; $\omega/\pi = 0.6$; $\alpha = \pi/16$



Final slide

Summary

- Attached masses excite resonance/anti-resonance to suppress low-frequency vibrations.
- Lightweight attachments can achieve this using geometry.

Next

- Nonlinearity (physical or geometrical) to transfer energy to higher frequencies (and damp).
- Modulational instability between acoustical and optical branches for efficient energy transfer.